

Federal University of Santa Catarina Graduate Program in Engineering and Mechanical Sciences

Plasmas and electrical discharges in gases (ECM410054)

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SUMMARY

Plasmas and electrical discharge in gases

- Kinetic theory of gases
- Atomic structure
- lonization
- Deionization
- Electron emission
- Behavior of charged particles in a gas in electric fields of low E/p
- Behavior of charged particles in a gas in electric fields of high E/p
- Glow discharges
- Plasmas



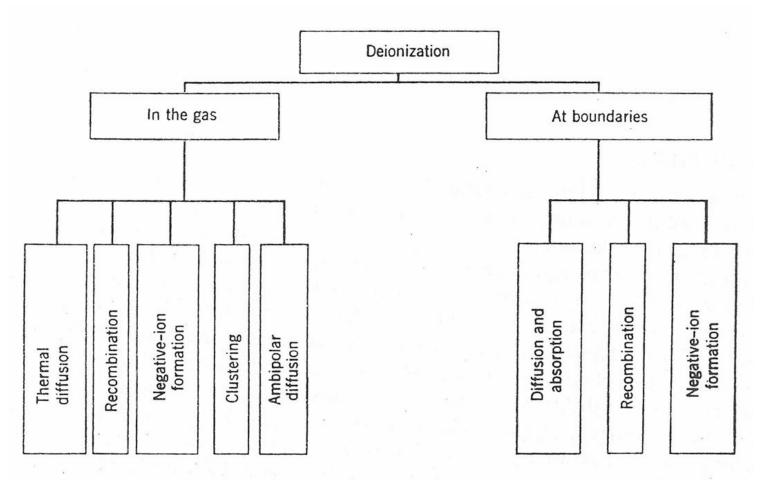
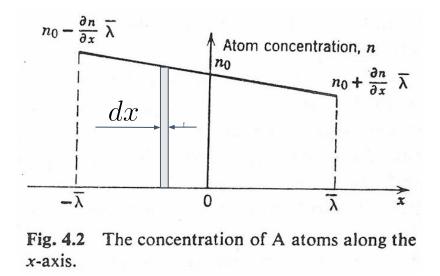


Fig. 4.1 Schematic diagram of possible deionization processes.



- Diffusion is caused by the difference in partial pressure or in molecular concentration. Consider a nonuniform distribution of gas A within a uniform gas B of higher concentration. Then, we shall state the following assumptions:
 - Number of atoms A per unit volume is much less than that of B.
 - All the A atoms have the same average speed v and the same mean free path λ .
 - Concentration of the A atoms changes in the *x*-direction only.



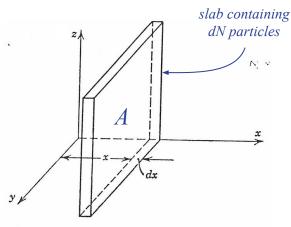


Fig. 4.4 A slab of unit cross section between x and x + dx for the evaluation of the change in particle density due to diffusion.



• Consider the case when the concentration gradient is linear:

$$n = n_0 + \left(\frac{\partial n}{\partial x}\right) x$$

• To calculate the number of A atoms crossing the plane perpendicular to the *x*-axis at x = 0 from the left to right and then vice-versa, we shall consider all the A atoms to the leftof this plane. Since motion in all directions is equally probable, then only one-sixth of the atoms will be moving in the positive x-direction:

$$n = \frac{dN}{dV} = \frac{dN}{Adx}$$
 : $\frac{dN}{A} = dj = ndx$

where dN is the number of particles in dx crossing the plane dA represented by dj. The total amount particles crossing the plane from left to the right j is:

$$j' = \frac{1}{6} \int_{-\lambda}^{0} \left[n_0 + \left(\frac{\partial n}{\partial x}\right) x \right] dx$$
$$j' = \frac{1}{6} \left[n_0 \lambda - \frac{1}{2} \left(\frac{\partial n}{\partial x}\right) \lambda^2 \right] = \frac{1}{6} \left[n_0 vt - \frac{1}{2} \left(\frac{\partial n}{\partial x}\right) (vt)^2 \right]$$



where t is the time between two successive collisions. The flow of charges is given by:

$$\frac{\partial j'}{\partial t} = \frac{1}{6} \left[n_0 v - \left(\frac{\partial n}{\partial x} \right) \lambda v \right]$$

• Similarly, we can calculate the number of A atoms crossing the plane x = 0 from right to left as:

$$j'' = \frac{1}{6} \int_{\lambda}^{0} \left[n_0 + \left(\frac{\partial n}{\partial x}\right) x \right] dx$$

$$\frac{\partial j''}{\partial t} = -\frac{1}{6} \left[n_0 v + \left(\frac{\partial n}{\partial x} \right) \lambda v \right]$$

• Then the resultant number of A atoms crossing the plane from left to the right is:

$$J = \frac{\partial j'}{\partial t} + \frac{\partial j''}{\partial t} = \frac{1}{6} \left[n_0 v - \left(\frac{\partial n}{\partial x}\right) \lambda v \right] - \frac{1}{6} \left[n_0 v + \left(\frac{\partial n}{\partial x}\right) \lambda v \right]$$



where *J* is the flow of charges (current density):

$$J = -\frac{1}{3}\lambda v \frac{\partial n}{\partial x}$$

• In three coordinates, the above equation is given by:

$$\vec{J} = -D\vec{\nabla}n$$

where $D = \lambda v/3$ is the diffusion coefficient. The above equation states the following assumptions:

- The negative sign indicates that the flow of atoms is opposite to the concentration gradient;
- The flow stops when the gradient is zero.



Considering the densities n⁺ and n⁻ of the positive and negative particles, respectively. The number of recombinations per unit time must be proportional to n⁺ and n⁻. If the constant of proportionality is ρ then:

$$dn = \boxed{-\rho n^+ n^- dt}$$

Radiative recombination

 $e^- + Ar^+ \rightarrow Ar + h\nu$

where here we assume the so-called quasi-neutrality condition $(n^+ = n^- = n)$. Then,

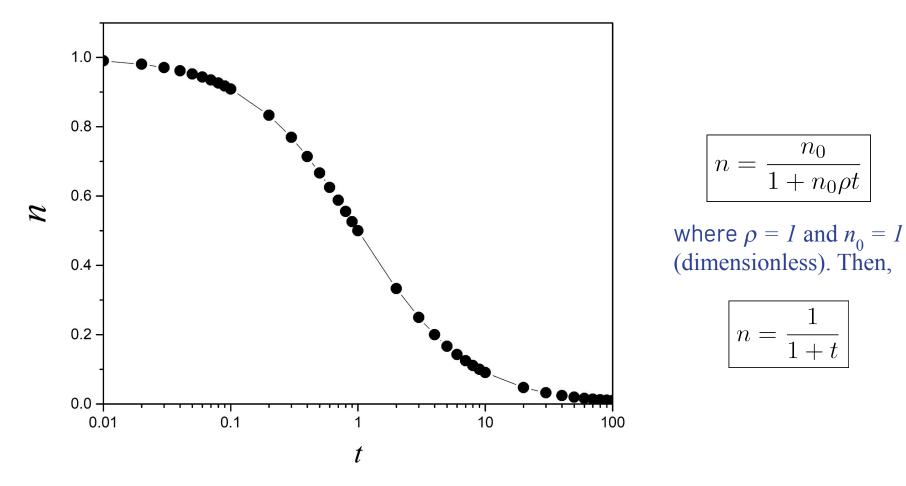
 $\begin{array}{c} & & & & & \\ & & & & \\ & & & & \\ & & & \\ & e^{-} & Ar^{+} & Ar \end{array} \end{array}$

$$dn = -\rho n^2 dt$$

$$\int_{n_0}^{n} \frac{dn}{n^2} = -\rho t \quad \therefore \quad \left| n = \frac{n_0}{1 + n_0 \rho t} \right|$$

that is defined as "ion-loss equation" where ρ is defined by plotting 1/n versus t.





 n_0

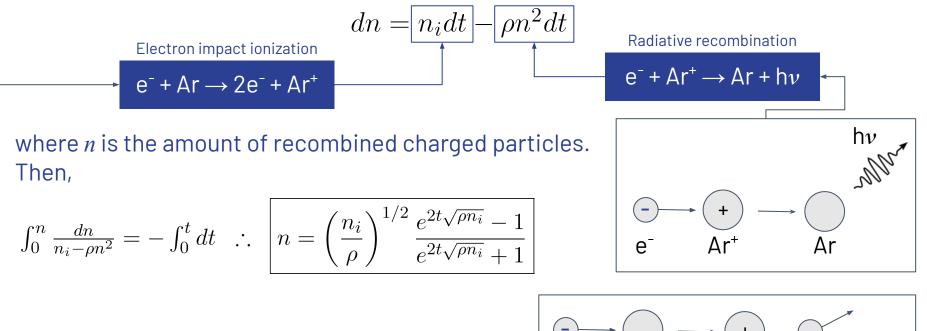
 $1 + n_0 \rho t$

 $n = \frac{1}{1+t}$

n =



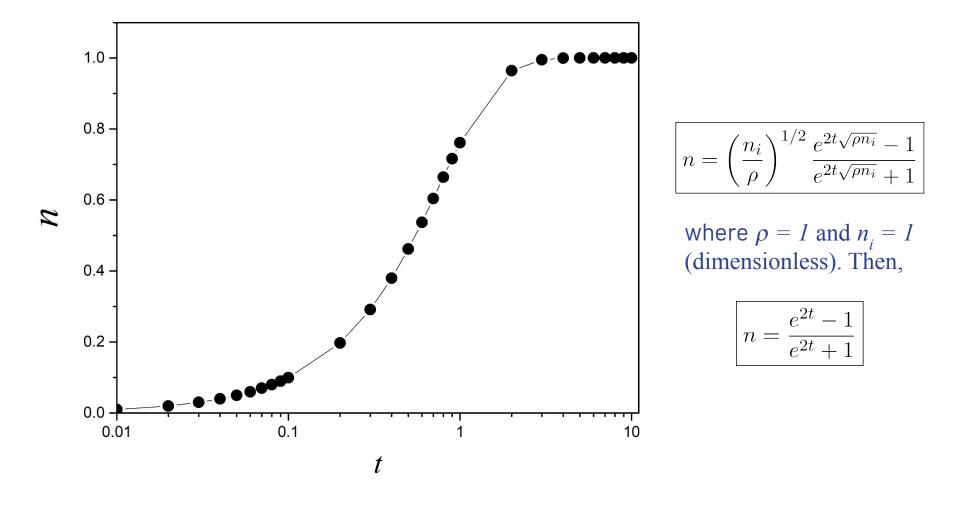
 If, however, the gas is subject to an external ionizer that continuously ionizes some neutral particles, an increase in the charge-particle density n will take place. If the density of newly ionized particles per second is n, we get:



e⁻

that is defined as "ionization-growth method".







DEIONIZATION

Recombination processes: electron attachment (examples)

• Radiative recombination:

 $e^- + Ar^+ \rightarrow Ar + h\nu$

• Dissociative recombination:

$$e^{-} + 0_{2}^{+} \rightarrow 0 + 0$$

• Dielectronic recombination:

 $e^- + Ar^* \rightarrow Ar^{**} (Ar^{**} \rightarrow Ar^+ + e^- (autoionization))$

• Three-body-recombination;

$$A^{+} + B + e^{-} \rightarrow A^{*} + B$$
$$A^{+} + 2e^{-} \rightarrow A^{*} + e^{-}$$



• Radiative attachment (the reverse case is called photodetachment):

 $e^- + 0 \rightarrow 0^- + hv$

• Three-body attachment:

 $e^- + A + B \rightarrow A^- + (B + W_K)$

• Dissociative attachment (the reverse case is called associative detachment):

 $e^- + AB \rightarrow A^- + B$

• Charge transfer between heavy particles:

$$A + B \rightarrow A^+ + B^-$$



• Chapter 4 - E. Nasser, Fundamentals of Gaseous Ionization and Plasma Electronics (pages 102-107; 114-118).

See you next topic!

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